

# Delta Function of a Complex Variable

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**Abstract** Ignorance of Infinitesimal Complex Calculus prevented the definition of the Complex Delta Function.

$\delta(\zeta - z)$  is a Plane Delta Function that spikes to  $\frac{1}{2\pi i(\zeta - z)}$  on the

Infinitesimal Disk  $|\zeta - z| \leq dr$ , and vanishes outside it.

For  $dr = \left\langle \frac{1}{n} \right\rangle$ ,  $r = |z|$ , and  $\phi = \text{Arg}(z)$

$$\delta(z) = \frac{1}{2\pi i} e^{-i\phi} \frac{1}{dr} e^{-\frac{r}{dr}}.$$

Delta has the Bessel Integral Representation

$$\delta(z) = \frac{1}{2\pi i} e^{-i\phi} r \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega r) \Omega d\Omega$$

The Circulation of Delta along the Infinitesimal circle  $|\zeta - z| = dr$

is

$$\oint_{|\zeta-z|=dr} \delta(\zeta - z) d\zeta = 1$$

And If  $f(z)$  is Hyper-Complex Differentiable function at  $z$

$$\text{Then, } \oint_{|\zeta-z|=dr} f(\zeta)\delta(\zeta-z)d\zeta = f(z).$$

For an analytic function on a simply connected domain, this Sifting through the function values by the Complex Delta Function is the Cauchy Integral Formula

Also the Residue of a Laurent Expansion of a singular Function follows from the sifting property of the Delta Function.

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# Contents

## Introduction

1. Hyper-real Line
2. Hyper-real Integral
3. Delta Function
4. The Fourier Transform
5. Hyper-Complex Plane
6. Hyper-Complex Path Integral
7. Hyper-Complex  $\delta(z)$
8.  $\delta(z)$  Plots
9.  $\delta(z)$  Properties
10.  $\delta(f(z))$
11. Polar Representation of  $\delta(z)$
12.  $\delta(\rho)$
13.  $\delta(\rho)$  and  $\delta(x)\delta(y)$
14. Cartesian Representation of  $\delta(z)$
15. Bessel Integral Representation of  $\delta(z)$
16. Primitive and Derivatives of  $\delta(z)$
17. Circulation of  $\delta(z)$
18. Sifting by  $\delta(\zeta - z)$ , and  $\frac{d}{dz}\delta(\zeta - z)$

19.  $\delta(\zeta - z)$  and the Cauchy Integral Formula

20.  $\delta(\zeta - z)$  and the Residue of a Laurent Expansion

References

# Introduction

## 0.1 Dirac's Confusion about the Complex Delta Function

The confusion about the Complex Delta Function can be traced to Dirac. In [Dirac] at the bottom of page 778, he writes

*"...,we must use the formula*

$$\frac{d}{dz} \log z = \frac{1}{z} - i\pi\delta(z), \quad (27)$$

*in which the term*

$$-i\pi\delta(z)$$

*is required in order to make the integral of the right-hand side of (27) between the limits  $a$ , and  $-a$  equal*

$$\log(-1),$$

*the integral of*

$$\frac{1}{z}$$

*between these limits being assumed to be zero."*

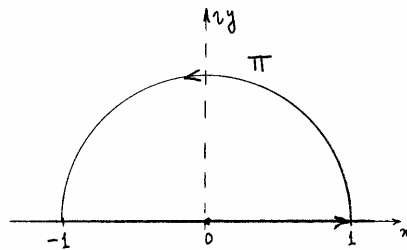
Dirac made here four errors:

**1<sup>st</sup> Error**,  $\log(-1) \neq -i\pi$

Indeed, integrating the left-hand side of (27), for  $a \neq 0$ ,

$$\int_{z=-a}^{z=a} \frac{d}{dz} \log z = \log z \Big|_{z=-a}^{z=a}$$

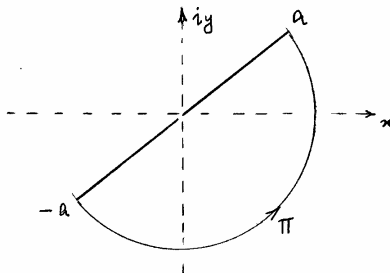
$$\begin{aligned}
 &= \log a - \log(-a) \\
 &= \log \frac{a}{-a} \\
 &= \log(-1) \\
 &= \underbrace{\log|-1|}_{=0} + i \underbrace{\text{Arg}(-1)}_{\pi} \\
 &= i\pi.
 \end{aligned}$$



**2<sup>nd</sup> Error**  $\int_{z=-a}^{z=a} \frac{1}{z} dz \neq 0$

Indeed, by the Residue Theorem

$$\begin{aligned}
 \int_{z=-a}^{z=a} \frac{1}{z} dz &= i \underbrace{[\text{Arg}(a) - \text{Arg}(-a)]}_{\pi} \\
 &= i\pi.
 \end{aligned}$$



$$\underline{\mathbf{3^{rd} Error} \quad \frac{d}{dz} \log z \neq \frac{1}{z} - i\pi\delta(z)}$$

Indeed, it is well known that without any exception

$$\frac{d}{dz} \log z = \frac{1}{z}.$$

$$\underline{\mathbf{4^{th} Error} \quad \int_{z=-a}^{z=a} \delta(z) dz \neq 1}$$

Indeed, unlike the hyper-real Delta Function  $\delta(x)$ , which sifts through its singularity at  $x = 0$ , the Hyper-Complex Delta Function sifts along a line that encircles the singularity, and avoids it. In the proceeding, it will become clear that

$$\begin{aligned} \int_{z=-a}^{z=a} \delta(z) dz &= \frac{1}{2\pi i} i \underbrace{[\text{Arg}(a) - \text{Arg}(-a)]}_{\pi} \\ &= \frac{1}{2}. \end{aligned}$$

Dirac's theory about the infinite distribution of electrons in the theory of the positron is flawed, because of Dirac's confusion about the Complex Delta Function.

## 0.2 The Hyper-Complex Delta Function

For  $z$  in the interior of  $\gamma$ , Cauchy Integral Formula gives an

analytic  $f(z)$  as the convolution of  $f$  with  $\frac{1}{2\pi i} \frac{1}{\zeta}$ .

$$f(z) = \oint_{\gamma} f(\zeta) \frac{1}{2\pi i} \frac{1}{\zeta - z} d\zeta$$

Thus,  $\frac{1}{2\pi i} \frac{1}{\zeta - z}$  recovers the value of a complex function  $f(\zeta)$  at the point  $z$  in the interior of a loop  $\gamma$ , by sifting through the values of  $f(\zeta)$  on  $\gamma$ .

However,  $\frac{1}{2\pi i} \frac{1}{\zeta - z}$  is the Hyper-Complex Delta Function only if the integration path is in the infinitesimal disk

$$|\zeta - z| \leq dr,$$

then, the sifting is performed by

$$\frac{1}{2\pi i} \frac{1}{\zeta - z} \mathcal{X}_{\{|\zeta - z| \leq dr\}}(\zeta) \equiv \delta(\zeta - z).$$

We call  $\delta(\zeta - z)$  the Hyper-Complex Delta Function.

In [Dan7], we showed that

$$\oint_{|\zeta - z| = dr} \frac{1}{2\pi i} \frac{1}{\zeta - z} d\zeta = 1.$$

Due to  $\frac{1}{\zeta - z}$ , the Complex Delta spikes to

$$\frac{1}{2\pi i} \frac{1}{d\zeta} = \frac{1}{2\pi i dr} e^{-i \text{Arg}(\zeta - z)}$$



on the infinitesimal disk  $|\zeta - z| \leq dr$ .

The primitive of the Complex Delta on the disk is

$$\frac{1}{2\pi i} \text{Log}(\zeta - z),$$

and the derivative of the Complex Delta on the disk is

$$\frac{1}{2\pi i} \frac{1}{(\zeta - z)^2}.$$

Since

$$\zeta - z = (dr)e^{-i \text{Arg}(\zeta - z)},$$

is a hyper-complex infinitesimal, we need to recall the Hyper-Complex Plane that we introduced in [Dan7].

Since the Complex Delta Function is an extension of the Real Delta Function, we need to recall the Delta Function that we introduced in [Dan4].

# 1.

## Hyper-real Line

Each real number  $\alpha$  can be represented by a Cauchy sequence of rational numbers,  $(r_1, r_2, r_3, \dots)$  so that  $r_n \rightarrow \alpha$ .

The constant sequence  $(\alpha, \alpha, \alpha, \dots)$  is a constant Hyper-real.

In [Dan2] we established that,

1. Any totally ordered set of positive, monotonically decreasing to zero sequences  $(l_1, l_2, l_3, \dots)$  constitutes a family of infinitesimal Hyper-reals.
2. The infinitesimals are smaller than any real number, yet strictly greater than zero.
3. Their reciprocals  $\left(\frac{1}{l_1}, \frac{1}{l_2}, \frac{1}{l_3}, \dots\right)$  are the infinite Hyper-reals.
4. The infinite Hyper-reals are greater than any real number, yet strictly smaller than infinity.
5. The infinite Hyper-reals with negative signs are smaller than any real number, yet strictly greater than  $-\infty$ .
6. The sum of a real number with an infinitesimal is a non-constant Hyper-real.

7. The Hyper-reals are the totality of constant Hyper-reals, a family of infinitesimals, a family of infinitesimals with negative sign, a family of infinite Hyper-reals, a family of infinite Hyper-reals with negative sign, and non-constant Hyper-reals.
8. The Hyper-reals are totally ordered, and aligned along a line: the Hyper-real Line.
9. That line includes the real numbers separated by the non-constant Hyper-reals. Each real number is the center of an interval of Hyper-reals, that includes no other real number.
10. In particular, zero is separated from any positive real by the infinitesimals, and from any negative real by the infinitesimals with negative signs,  $-dx$ .
11. Zero is not an infinitesimal, because zero is not strictly greater than zero.
12. We do not add infinity to the Hyper-real line.
13. The infinitesimals, the infinitesimals with negative signs, the infinite Hyper-reals, and the infinite Hyper-reals with negative signs are semi-groups with respect to addition. Neither set includes zero.

14. The Hyper-real line is embedded in  $\mathbb{R}^\infty$ , and is not homeomorphic to the real line. There is no bi-continuous one-one mapping from the Hyper-real onto the real line.
15. In particular, there are no points on the real line that can be assigned uniquely to the infinitesimal Hyper-reals, or to the infinite Hyper-reals, or to the non-constant Hyper-reals.
16. No neighbourhood of a Hyper-real is homeomorphic to an  $\mathbb{R}^n$  ball. Therefore, the Hyper-real line is not a manifold.
17. The Hyper-real line is totally ordered like a line, but it is not spanned by one element, and it is not one-dimensional.

## 2.

# Hyper-real Integral

In [Dan3], we defined the integral of a Hyper-real Function.

Let  $f(x)$  be a Hyper-real function on the interval  $[a, b]$ .

The interval may not be bounded.

$f(x)$  may take infinite Hyper-real values, and need not be bounded.

At each

$$a \leq x \leq b,$$

there is a rectangle with base  $[x - \frac{dx}{2}, x + \frac{dx}{2}]$ , height  $f(x)$ , and area

$$f(x)dx.$$

We form the **Integration Sum** of all the areas for the  $x$ 's that start at  $x = a$ , and end at  $x = b$ ,

$$\sum_{x \in [a, b]} f(x)dx.$$

If for any infinitesimal  $dx$ , the Integration Sum has the same Hyper-real value, then  $f(x)$  is integrable over the interval  $[a, b]$ .

Then, we call the Integration Sum the integral of  $f(x)$  from  $x = a$ , to  $x = b$ , and denote it by

$$\int_{x=a}^{x=b} f(x)dx .$$

If the Hyper-real is infinite, then it is the integral over  $[a, b]$ ,

If the Hyper-real is finite,

$$\int_{x=a}^{x=b} f(x)dx = \text{real part of the hyper-real} . \square$$

## 2.1 The countability of the Integration Sum

In [Dan1], we established the equality of all positive infinities:

We proved that the number of the Natural Numbers,

$Card\mathbb{N}$ , equals the number of Real Numbers,  $Card\mathbb{R} = 2^{Card\mathbb{N}}$ , and

we have

$$Card\mathbb{N} = (Card\mathbb{N})^2 = \dots = 2^{Card\mathbb{N}} = 2^{2^{Card\mathbb{N}}} = \dots \equiv \infty .$$

In particular, we demonstrated that the real numbers may be well-ordered.

Consequently, there are countably many real numbers in the interval  $[a, b]$ , and the Integration Sum has countably many terms.

While we do not sequence the real numbers in the interval, the summation takes place over countably many  $f(x)dx$ .

The Lower Integral is the Integration Sum where  $f(x)$  is replaced by its lowest value on each interval  $[x - \frac{dx}{2}, x + \frac{dx}{2}]$

$$\mathbf{2.2} \quad \sum_{x \in [a, b]} \left( \inf_{x - \frac{dx}{2} \leq t \leq x + \frac{dx}{2}} f(t) \right) dx$$

The Upper Integral is the Integration Sum where  $f(x)$  is replaced by its largest value on each interval  $[x - \frac{dx}{2}, x + \frac{dx}{2}]$

$$\mathbf{2.3} \quad \sum_{x \in [a, b]} \left( \sup_{x - \frac{dx}{2} \leq t \leq x + \frac{dx}{2}} f(t) \right) dx$$

If the integral is a finite Hyper-real, we have

**2.4** *A Hyper-real function has a finite integral if and only if its upper integral and its lower integral are finite, and differ by an infinitesimal.*

### 3.

## Delta Function

In [Dan5], we defined the Delta Function, and established its properties

1. The Delta Function is a Hyper-real function defined from the Hyper-real line into the set of two Hyper-reals  $\left\{0, \frac{1}{dx}\right\}$ . The Hyper-real 0 is the sequence  $\langle 0, 0, 0, \dots \rangle$ . The infinite Hyper-real  $\frac{1}{dx}$  depends on our choice of  $dx$ .
2. We will usually choose the family of infinitesimals that is spanned by the sequences  $\left\langle \frac{1}{n} \right\rangle, \left\langle \frac{1}{n^2} \right\rangle, \left\langle \frac{1}{n^3} \right\rangle, \dots$ . It is a semigroup with respect to vector addition, and includes all the scalar multiples of the generating sequences that are non-zero. That is, the family includes infinitesimals with negative sign. Therefore,  $\frac{1}{dx}$  will mean the sequence  $\langle n \rangle$ . Alternatively, we may choose the family spanned by the



sequences  $\left\langle \frac{1}{2^n} \right\rangle, \left\langle \frac{1}{3^n} \right\rangle, \left\langle \frac{1}{4^n} \right\rangle, \dots$ . Then,  $\frac{1}{dx}$  will mean the sequence  $\left\langle 2^n \right\rangle$ . Once we determined the basic infinitesimal  $dx$ , we will use it in the Infinite Riemann Sum that defines an Integral in Infinitesimal Calculus.

3. The Delta Function is strictly smaller than  $\infty$

4. We define,  $\delta(x) \equiv \frac{1}{dx} \chi_{\left[-\frac{dx}{2}, \frac{dx}{2}\right]}(x)$ ,

$$\text{where } \chi_{\left[-\frac{dx}{2}, \frac{dx}{2}\right]}(x) = \begin{cases} 1, & x \in \left[-\frac{dx}{2}, \frac{dx}{2}\right] \\ 0, & \text{otherwise} \end{cases}.$$

5. Hence,

$$\diamond \text{ for } x < 0, \delta(x) = 0$$

$$\diamond \text{ at } x = -\frac{dx}{2}, \delta(x) \text{ jumps from } 0 \text{ to } \frac{1}{dx},$$

$$\diamond \text{ for } x \in \left[-\frac{dx}{2}, \frac{dx}{2}\right], \delta(x) = \frac{1}{dx}.$$

$$\diamond \text{ at } x = 0, \delta(0) = \frac{1}{dx}$$

$$\diamond \text{ at } x = \frac{dx}{2}, \delta(x) \text{ drops from } \frac{1}{dx} \text{ to } 0.$$

$$\diamond \text{ for } x > 0, \delta(x) = 0.$$

$$\diamond x\delta(x) = 0$$

6. If  $dx = \langle \frac{1}{n} \rangle$ ,  $\delta(x) = \langle \chi_{[-\frac{1}{2}, \frac{1}{2}]}(x), 2\chi_{[-\frac{1}{4}, \frac{1}{4}]}(x), 3\chi_{[-\frac{1}{6}, \frac{1}{6}]}(x) \dots \rangle$

7. If  $dx = \langle \frac{2}{n} \rangle$ ,  $\delta(x) = \langle \frac{1}{2 \cosh^2 x}, \frac{2}{2 \cosh^2 2x}, \frac{3}{2 \cosh^2 3x}, \dots \rangle$

8. If  $dx = \langle \frac{1}{n} \rangle$ ,  $\delta(x) = \langle e^{-x} \chi_{[0, \infty)}, 2e^{-2x} \chi_{[0, \infty)}, 3e^{-3x} \chi_{[0, \infty)}, \dots \rangle$

9.  $\int_{x=-\infty}^{x=\infty} \delta(x) dx = 1.$

## 4.

# The Fourier Transform

In [Dan6], we defined the Fourier Transform and established its properties

1.  $\mathcal{F}\{\delta(x)\} = 1$
2.  $\delta(x) = \text{the inverse Fourier Transform of the unit function } 1$

$$= \frac{1}{2\pi} \int_{\omega=-\infty}^{\omega=\infty} e^{i\omega x} d\omega$$

$$= \int_{\nu=-\infty}^{\nu=\infty} e^{2\pi i x} d\nu, \quad \omega = 2\pi\nu$$

$$3. \left. \frac{1}{2\pi} \int_{\omega=-\infty}^{\omega=\infty} e^{i\omega x} d\omega \right|_{x=0} = \frac{1}{dx} = \text{an infinite Hyper-real}$$

$$\left. \int_{\omega=-\infty}^{\omega=\infty} e^{i\omega x} d\omega \right|_{x \neq 0} = 0$$

### 4. Fourier Integral Theorem

$$f(x) = \frac{1}{2\pi} \int_{k=-\infty}^{k=\infty} \left( \int_{\xi=-\infty}^{\xi=\infty} f(\xi) e^{-ik\xi} d\xi \right) e^{ikx} dk$$

does not hold in the Calculus of Limits, under any

conditions.

## 5. Fourier Integral Theorem in Infinitesimal Calculus

If  $f(x)$  is a Hyper-real function,

Then,

➤ the Fourier Integral Theorem holds.

$$\text{➤ } \int_{x=-\infty}^{x=\infty} f(x)e^{-i\alpha x} dx \text{ converges to } F(\alpha)$$

$$\text{➤ } \frac{1}{2\pi} \int_{\alpha=-\infty}^{\alpha=\infty} F(\alpha)e^{-i\alpha x} d\alpha \text{ converges to } f(x)$$

## 6. 2-Dimensional Fourier Transform

$$\begin{aligned} \mathcal{F}\{f(x, y)\} &= \int_{y=-\infty}^{y=\infty} \int_{x=-\infty}^{x=\infty} f(x, y)e^{-i\omega_x x - i\omega_y y} dx dy \\ &= \int_{y=-\infty}^{y=\infty} \int_{x=-\infty}^{x=\infty} f(x, y)e^{-2\pi i(\nu_x x + \nu_y y)} dx dy, \quad \begin{aligned} \omega_x &= 2\pi\nu_x \\ \omega_y &= 2\pi\nu_y \end{aligned} \end{aligned}$$

## 8. 2-Dimensional Inverse Fourier Transform

$$\mathcal{F}^{-1}\{F(\omega_x, \omega_y)\} = \frac{1}{(2\pi)^2} \int_{\omega_y=-\infty}^{\omega_y=\infty} \int_{\omega_x=-\infty}^{\omega_x=\infty} F(\omega_x, \omega_y)e^{i(\omega_x x + \omega_y y)} d\omega_x d\omega_y$$

$$= \int_{\nu_y=-\infty}^{\nu_y=\infty} \int_{\nu_x=-\infty}^{\nu_x=\infty} F(2\pi\nu_x, 2\pi\nu_y) e^{2\pi i(\nu_x x + \nu_y y)} d\nu_x d\nu_y, \quad \begin{aligned} \omega_x &= 2\pi\nu_x \\ \omega_y &= 2\pi\nu_y \end{aligned}$$

## 9. 2-Dimensional Fourier Integral Theorem

$$\begin{aligned} f(x, y) &= \frac{1}{(2\pi)^2} \int_{\omega_y=-\infty}^{\omega_y=\infty} \int_{\omega_x=-\infty}^{\omega_x=\infty} \left( \int_{\eta=-\infty}^{\eta=\infty} \int_{\xi=-\infty}^{\xi=\infty} f(\xi, \eta) e^{-i\omega_x \xi - i\omega_y \eta} d\xi d\eta \right) e^{i(\omega_x x + \omega_y y)} d\omega_x d\omega_y \\ &= \int_{\eta=-\infty}^{\eta=\infty} \int_{\xi=-\infty}^{\xi=\infty} f(\xi, \eta) \left( \frac{1}{2\pi} \int_{\omega_x=-\infty}^{\omega_x=\infty} e^{i\omega_x(x-\xi)} d\omega_x \right) d\xi \left( \frac{1}{2\pi} \int_{\omega_y=-\infty}^{\omega_y=\infty} e^{i\omega_y(y-\eta)} d\omega_y \right) d\eta \\ &= \int_{\eta=-\infty}^{\eta=\infty} \int_{\xi=-\infty}^{\xi=\infty} f(\xi, \eta) \left( \int_{\nu_x=-\infty}^{\nu_x=\infty} e^{2\pi i \nu_x(x-\xi)} d\nu_x \right) d\xi \left( \int_{\nu_y=-\infty}^{\nu_y=\infty} e^{2\pi i \nu_y(y-\eta)} d\nu_y \right) d\eta, \quad \begin{aligned} \omega_x &= 2\pi\nu_x \\ \omega_y &= 2\pi\nu_y \end{aligned} \end{aligned}$$

## 10. 2-Dimensional Delta Function

$$\begin{aligned} \delta(x, y) &= \left( \frac{1}{2\pi} \int_{\omega_x=-\infty}^{\omega_x=\infty} e^{i\omega_x x} d\omega_x \right) \left( \frac{1}{2\pi} \int_{\omega_y=-\infty}^{\omega_y=\infty} e^{i\omega_y y} d\omega_y \right) \\ &= \left( \int_{\nu_x=-\infty}^{\nu_x=\infty} e^{2\pi i \nu_x x} d\nu_x \right) \left( \int_{\nu_y=-\infty}^{\nu_y=\infty} e^{2\pi i \nu_y y} d\nu_y \right), \quad \begin{aligned} \omega_x &= 2\pi\nu_x \\ \omega_y &= 2\pi\nu_y \end{aligned} \end{aligned}$$

## 5.

# Hyper-Complex Plane

Each complex number  $\alpha + i\beta$  can be represented by a Cauchy sequence of rational complex numbers,  $\langle r_1 + is_1, r_2 + is_2, r_3 + is_3, \dots \rangle$  so that  $r_n + is_n \rightarrow \alpha + i\beta$ .

The constant sequence  $(\alpha + i\beta, \alpha + i\beta, \alpha + i\beta, \dots)$  is a Constant Hyper-Complex Number.

In [Dan2] we established that,

1. Any set of sequences  $(l_1 + io_1, l_2 + io_2, l_3 + io_3, \dots)$ , where  $(l_1, l_2, l_3, \dots)$  belongs to one family of infinitesimal hyper reals, and  $(o_1, o_2, o_3, \dots)$  belongs to another family of infinitesimal hyper-reals, constitutes a family of infinitesimal hyper-complex numbers.
2. Each hyper-complex infinitesimal has a polar representation  $dz = (dr)e^{i\phi} = o_*e^{i\phi}$ , where  $dr = o_*$  is an infinitesimal, and  $\phi = \arg(dz)$ .
3. The infinitesimal hyper-complex numbers are smaller in length, than any complex number, yet strictly greater than

zero.

4. Their reciprocals  $\left(\frac{1}{t_1+io_1}, \frac{1}{t_2+io_2}, \frac{1}{t_3+io_3}, \dots\right)$  are the infinite hyper-complex numbers.
5. The infinite hyper-complex numbers are greater in length than any complex number, yet strictly smaller than infinity.
6. The sum of a complex number with an infinitesimal hyper-complex is a non-constant hyper-complex.
7. The Hyper-Complex Numbers are the totality of constant hyper-complex numbers, a family of hyper-complex infinitesimals, a family of infinite hyper-complex, and non-constant hyper-complex.
8. The Hyper-Complex Plane is the direct product of a Hyper-Real Line by an imaginary Hyper-Real Line.
9. In Cartesian Coordinates, the Hyper-Real Line serves as an  $x$  coordinate line, and the imaginary as an  $iy$  coordinate line.
10. In Polar Coordinates, the Hyper-Real Line serves as a Range  $r$  line, and the imaginary as an  $i\theta$  coordinate. Radial symmetry leads to Polar Coordinates.
11. The Hyper-Complex Plane includes the complex numbers separated by the non-constant hyper-complex

- numbers. Each complex number is the center of a disk of hyper-complex numbers, that includes no other complex number.
12. In particular, zero is separated from any complex number by a disk of complex infinitesimals.
  13. Zero is not a complex infinitesimal, because the length of zero is not strictly greater than zero.
  14. We do not add infinity to the hyper-complex plane.
  15. The hyper-complex plane is embedded in  $\mathbb{C}^\infty$ , and is not homeomorphic to the Complex Plane  $\mathbb{C}$ . There is no bi-continuous one-one mapping from the hyper-complex Plane onto the Complex Plane.
  16. In particular, there are no points in the Complex Plane that can be assigned uniquely to the hyper-complex infinitesimals, or to the infinite hyper-complex numbers, or to the non-constant hyper-complex numbers.
  17. No neighbourhood of a hyper-complex number is homeomorphic to a  $\mathbb{C}^n$  ball. Therefore, the Hyper-Complex Plane is not a manifold.
  18. The Hyper-Complex Plane is not spanned by two elements, and is not two-dimensional.



## 6.

# Hyper-Complex Path Integral

Following the definition of the Hyper-real Integral in [Dan3], the Hyper-Complex Integral of  $f(z)$  over a path  $z(t)$ ,  $t \in [\alpha, \beta]$ , in its domain, is the sum of the areas  $f(z)z'(t)dt = f(z)dz(t)$  of the rectangles with base  $z'(t)dt = dz$ , and height  $f(z)$ .

### 6.1 Hyper-Complex Path Integral Definition

Let  $f(z)$  be hyper-complex function, defined on a domain in the Hyper-Complex Plane. The domain may not be bounded.

$f(z)$  may take infinite hyper-complex values, and need not be bounded.

Let  $z(t)$ ,  $t \in [\alpha, \beta]$ , be a path,  $\gamma(a, b)$ , so that  $dz = z'(t)dt$ , and  $z'(t)$  is continuous.

For each  $t$ , there is a hyper-complex rectangle with base  $[z(t) - \frac{dz}{2}, z(t) + \frac{dz}{2}]$ , height  $f(z)$ , and area  $f(z(t))dz(t)$ .

We form the **Integration Sum** of all the areas that start at  $z(\alpha) = a$ , and end at  $z(\beta) = b$ ,

$$\sum_{t \in [\alpha, \beta]} f(z(t)) dz(t).$$

If for any infinitesimal  $dz = z'(t)dt$ , the Integration Sum equals the same hyper-complex number, then  $f(z)$  is Hyper-Complex Integrable over the path  $\gamma(a, b)$ .

Then, we call the Integration Sum the Hyper-Complex Integral of  $f(z)$  over the  $\gamma(a, b)$ , and denote it by  $\int_{\gamma(a, b)} f(z) dz$ .

If the hyper-complex number is an infinite hyper-complex, then it equals  $\int_{\gamma(a, b)} f(z) dz$ .

If the hyper-complex number is finite, then its constant part equals  $\int_{\gamma(a, b)} f(z) dz$ .  $\square$

The Integration Sum may take infinite hyper-complex values, such as  $\frac{1}{dz}$ , but may not equal to  $\infty$ .

The Hyper-Complex Integral of the function  $f(z) = \frac{1}{|z|}$  over a path that goes through  $z = 0$  diverges.

## 6.2 The Countability of the Integration Sum

In [Dan1], we established the equality of all positive infinities:

We proved that the number of the Natural Numbers,

$Card\mathbb{N}$ , equals the number of Real Numbers,  $Card\mathbb{R} = 2^{Card\mathbb{N}}$ , and we have

$$Card\mathbb{N} = (Card\mathbb{N})^2 = \dots = 2^{Card\mathbb{N}} = 2^{2^{Card\mathbb{N}}} = \dots \equiv \infty.$$

In particular, we demonstrated that the real numbers may be well-ordered.

Consequently, there are countably many real numbers in the interval  $[\alpha, \beta]$ , and the Integration Sum has countably many terms.

While we do not sequence the real numbers in the interval, the summation takes place over countably many  $f(z)dz$ .

### 6.3 Continuous $f(z)$ is Path-Integrable

*Hyper-Complex  $f(z)$  Continuous on  $D$  is Path-Integrable on  $D$*

*Proof:*

Let  $z(t)$ ,  $t \in [\alpha, \beta]$ , be a path,  $\gamma(a, b)$ , so that  $dz = z'(t)dt$ , and  $z'(t)$  is continuous. Then,

$$f(z(t))z'(t) = (u(x(t), y(t)) + iv(x(t), y(t)))(x'(t) + iy'(t))$$

$$\begin{aligned}
&= \underbrace{\left[ u(x(t), y(t))x'(t) - v(x(t), y(t))y'(t) \right]}_{U(t)} + \\
&\quad + i \underbrace{\left[ u(x(t), y(t))y'(t) + v(x(t), y(t))x'(t) \right]}_{V(t)} \\
&= U(t) + iV(t),
\end{aligned}$$

where  $U(t)$ , and  $V(t)$  are Hyper-Real Continuous on  $[\alpha, \beta]$ .

Therefore, by [Dan3, 12.4],  $U(t)$ , and  $V(t)$  are integrable on  $[\alpha, \beta]$ .

Hence,  $f(z(t))z'(t)$  is integrable on  $[\alpha, \beta]$ .

Since

$$\int_{t=\alpha}^{t=\beta} f(z(t))z'(t)dt = \int_{\gamma(a,b)} f(z)dz,$$

$f(z)$  is Path-Integrable on  $\gamma(a, b)$ .  $\square$

## 7.

# Hyper-Complex Delta $\delta(z)$

### 7.1 Domain and Range of $\delta(z)$

The Hyper-Complex Delta Function is defined from the Hyper-Complex plane into the set of two hyper-complex numbers,

$$\left\{ 0, \frac{1}{2\pi i dz} \right\}.$$

The hyper-complex 0 is the sequence  $\langle 0, 0, 0, \dots \rangle$ .

The infinite hyper-complex

$$\frac{1}{2\pi i} \frac{1}{dz} = \frac{1}{2\pi i} \frac{1}{dr} e^{-i\phi}$$

depends on

$$\text{Arg } z = \phi,$$

and on our choice of the infinitesimal  $dr$ .

We will usually choose the family of infinitesimals that is spanned

by the sequences  $\left\langle \frac{1}{n} \right\rangle, \left\langle \frac{1}{n^2} \right\rangle, \left\langle \frac{1}{n^3} \right\rangle, \dots$ . It is a semigroup with

respect to vector addition, and includes all the scalar multiples of the generating sequences that are non-zero. That is, the family includes infinitesimals with negative sign.

Therefore,  $\frac{1}{dr}$  will mean the sequence  $\langle n \rangle$ .

Alternatively, we may choose the family spanned by the sequences

$\left\langle \frac{1}{2^n} \right\rangle, \left\langle \frac{1}{2^{2n}} \right\rangle, \left\langle \frac{1}{2^{3n}} \right\rangle, \dots$ . Then,  $\frac{1}{dr}$  will mean the sequence  $\langle 2^n \rangle$ .

**7.2**  $\delta(z)$  is an infinite hyper-complex function on the infinitesimal

hyper-complex disk  $|z| \leq dr$ , In particular,

$$\boxed{\delta(z) < \infty}$$

*Proof:* Since  $dr > 0$ , we have  $\frac{1}{dr} < \infty$ , and  $\frac{1}{2\pi i dr} e^{-i\phi} < \infty$ .

**7.3 Definition of  $\delta(z - z_0)$**

For any  $dz$ ,

$$\begin{aligned} \delta(z - z_0) &\equiv \frac{1}{2\pi i dz} \chi_{\{|z-z_0| \leq |dz|\}}(z), \\ &= \frac{1}{2\pi i(dx + idy)} \begin{cases} 1, & |z - z_0| \leq \sqrt{(dx)^2 + (dy)^2} \\ 0, & \text{otherwise} \end{cases} \end{aligned}$$

Thus,

- on the disk,  $|z - z_0| \leq dr$ ,  $\delta(z - z_0) = \frac{1}{2\pi i} \frac{1}{dz}$ .
- off the disk, for  $|z - z_0| > dr$ ,  $\delta(z - z_0) = 0$ .

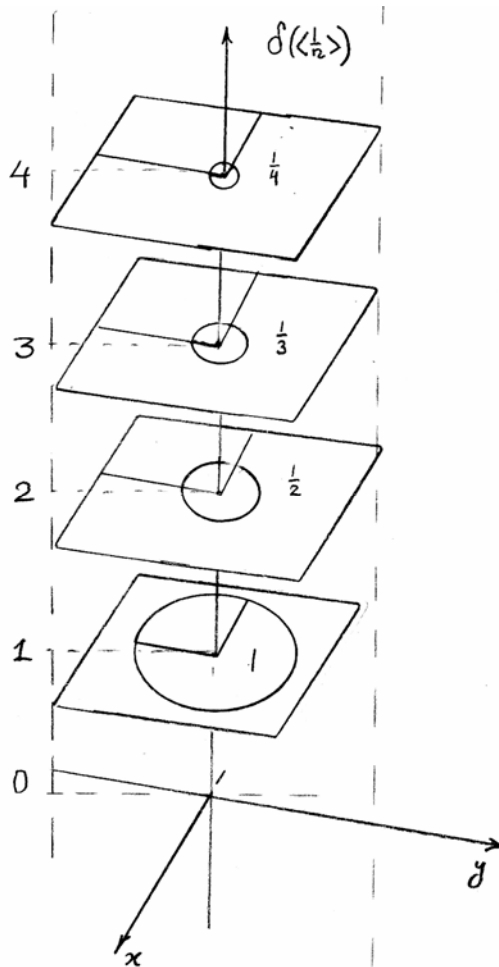
# 8.

## $\delta(z)$ Plots

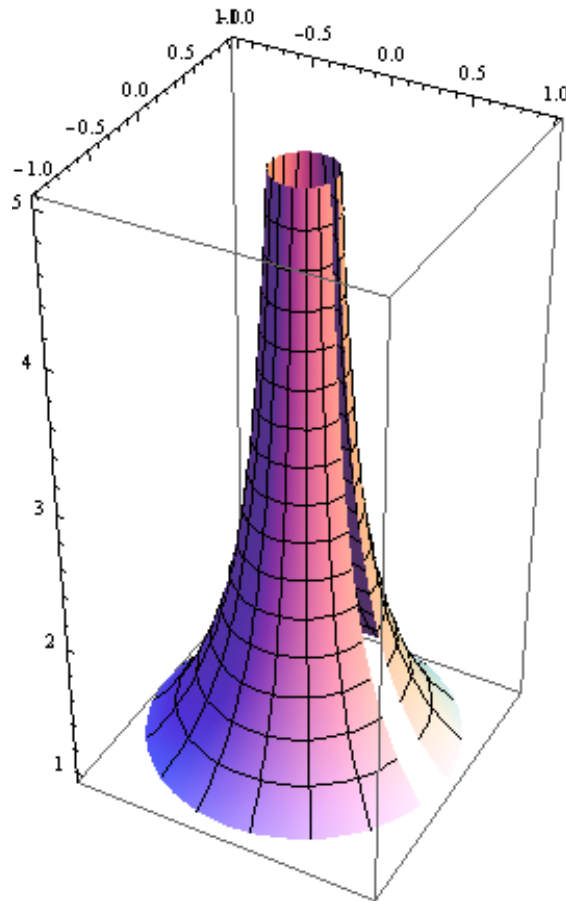
### 8.1 Delta Plot for

$$\delta(z) = \frac{1}{2\pi i} e^{-i\phi} \left\langle \chi_{\{|z|\leq 1\}}(z), 2\chi_{\{|z|\leq \frac{1}{2}\}}(z), 3\chi_{\{|z|\leq \frac{1}{3}\}}(z), \dots \right\rangle,$$

is a sequence of shrinking disks with increasing heights



The disks are in the body generated by revolving  $\frac{1}{|z|} = \frac{1}{r}$ , over the  
 Complex plane



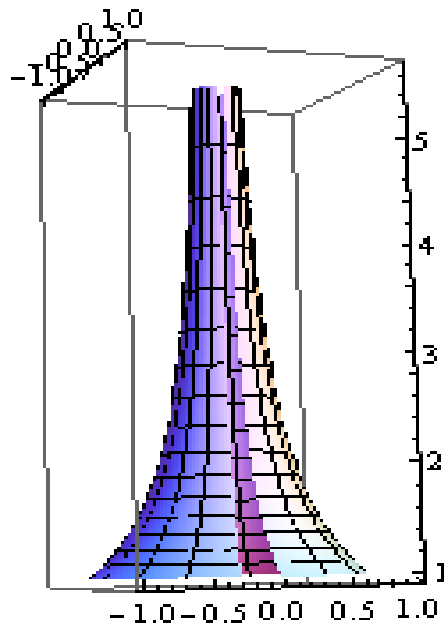
### 8.2 Delta Plot for

$$\delta(z) = \frac{1}{2\pi i} e^{-i\phi} \left\langle \chi_{\{|z| \leq \frac{1}{2^0}\}}(z), 2\chi_{\{|z| \leq \frac{1}{2^1}\}}(z), 4\chi_{\{|z| \leq \frac{1}{2^2}\}}(z), 8\chi_{\{|z| \leq \frac{1}{2^3}\}}(z) \dots \right\rangle$$



is a sequence of shrinking disks with increasing heights that lie in

the body of revolution of  $\frac{1}{|2^z|} = \frac{1}{2^x}$ .



**9.** **$\delta(z)$  Properties**

$$\mathbf{9.1} \quad \boxed{\delta(z - z_0) = \frac{1}{2\pi i} \frac{1}{dr} e^{-i(\phi - \phi_0)} \mathcal{X}_{\{|z - z_0| \leq dr\}}(z)},$$

$$\phi = \arg z, \phi_0 = \arg z_0$$

$$\mathbf{9.2} \quad \boxed{\delta(0) = \frac{1}{2\pi i dr}}$$

$$\mathbf{9.3} \quad \boxed{(\delta(z))^n = \frac{1}{(2\pi i)^n} \frac{1}{(dr)^n} e^{-in\phi} \mathcal{X}_{\{|z| \leq dr\}}(z)}, \quad n = 2, 3, \dots$$

**10.**

$$\delta(f(z))$$

$$\mathbf{10.1} \quad \boxed{\delta(az) = \frac{1}{a} \delta(z)}$$

*Proof:* 
$$\begin{aligned} \delta(az) &= \frac{1}{2\pi i} \frac{1}{d(az)} \mathcal{X}_{\{|az| \leq dr\}}(z) \\ &= \frac{1}{a} \frac{1}{2\pi i} \frac{1}{dz} \mathcal{X}_{\{|z| \leq \frac{1}{|a|} dr\}}(z) \\ &= \frac{1}{a} \frac{1}{2\pi i} \frac{1}{dz} \mathcal{X}_{\{|z| \leq |dz|\}}(z) \\ &= \frac{1}{a} \frac{1}{2\pi i} \delta(z). \quad \square \end{aligned}$$

$$\mathbf{10.2} \quad z_1 = \text{only zero of } f(z), \quad f'(z_1) \neq 0 \Rightarrow$$

$$\Rightarrow \boxed{\delta(f(z)) = \frac{1}{f'(z_1)} \delta(z - z_1)}$$

*Proof:* 
$$\delta(f(z)) = \delta(f(z) - f(z_1))$$

For  $z - z_1 = \text{infinitesimal}$ ,

$$\begin{aligned}
&= \delta(f'(z_1)(z - z_1)) \\
&= \frac{1}{f'(z_1)} \delta(z - z_1), \text{ by 10.1. } \square
\end{aligned}$$

**10.3**  $z_1, z_2$  are the only zeros of  $f(z); f'(z_1), f'(z_2) \neq 0 \Rightarrow$

$$\Rightarrow \boxed{\delta(f(z)) = \frac{1}{f'(z_1)} \delta(z - z_1) + \frac{1}{f'(z_2)} \delta(z - z_2)}$$

*Proof:*  $\delta(f(z)) = \delta(f(z) - f(z_1)) + \delta(f(z) - f(z_2))$

**For**  $z - z_1 = \text{infinitesimal}, z - z_2 = \text{infinitesimal},$

$$\begin{aligned}
\delta(f(z)) &= \delta(f'(z_1)(z - z_1)) + \delta(f'(z_2)(z - z_2)) \\
&= \frac{1}{f'(z_1)} \delta(z - z_1) + \frac{1}{f'(z_2)} \delta(z - z_2). \square
\end{aligned}$$

**10.4**  $\boxed{\delta(z^2 - a^2) = \frac{1}{2a} \delta(z - a) + \frac{1}{2a} \delta(z + a)}$

**10.5**  $\boxed{\delta((z - a)(z - b)) = \frac{1}{a - b} \delta(z - a) + \frac{1}{b - a} \delta(z - b)}$

**10.6**  $z_1, \dots, z_n$  are the only zeros of  $f(z)$ ;  $f'(z_1), \dots, f'(z_n) \neq 0 \Rightarrow$

$$\delta(f(z)) = \frac{1}{f'(z_1)} \delta(z - z_1) + \dots + \frac{1}{f'(z_n)} \delta(z - z_n)$$

**10.7**

$$\delta(\sin z) = \dots + \delta(z + 2\pi) - \delta(z + \pi) + \delta(z) - \delta(z - \pi) + \delta(z - 2\pi) + \dots$$

*Proof:* The zeros of  $\sin z$  are  $\dots - 2\pi, -\pi, 0, \pi, 2\pi, \dots$

and  $\dots, \cos(-\pi) = -1, \cos(0) = 1, \cos(\pi) = -1, \dots \square$

# 11.

## Polar Representation of $\delta(z)$

**11.1** 
$$\delta(z) = \frac{1}{2\pi i} e^{-i \text{Arg}(z)} \delta(|z|)$$

Proof:

$$\begin{aligned} \delta(z) &= \frac{1}{2\pi i dz} \mathcal{X}_{\{|\zeta| \leq |dz|\}}(z), \quad \text{where } \mathcal{X}_{\{|\zeta| \leq |dz|\}}(z) = \begin{cases} 0, & |\zeta| > |dz| \\ 1, & |\zeta| \leq |dz| \end{cases} \\ &= \frac{dx - idy}{2\pi i [(dx)^2 + (dy)^2]} \mathcal{X}_{\{|\zeta| \leq \sqrt{(dx)^2 + (dy)^2}\}}(z) \\ &= \frac{\frac{dx}{\sqrt{(dx)^2 + (dy)^2}} - i \frac{dy}{\sqrt{(dx)^2 + (dy)^2}}}{2\pi i \sqrt{(dx)^2 + (dy)^2}} \mathcal{X}_{\{|\zeta| \leq \sqrt{(dx)^2 + (dy)^2}\}}(z) \\ &= \frac{\cos \phi - i \sin \phi}{2\pi i} \underbrace{\frac{1}{d\rho} \mathcal{X}_{\{|\zeta| \leq d\rho\}}(z)}_{\delta(\rho)}, \quad \text{where } \begin{cases} d\rho = \sqrt{(dx)^2 + (dy)^2} \\ \phi = \text{Arg}(z) \end{cases} \\ &= \frac{1}{2\pi i} e^{-i \text{Arg}(z)} \delta(\rho), \quad \rho = |z| = |x + iy| = \sqrt{x^2 + y^2}. \end{aligned}$$

## 12.

$\delta(\rho)$

**12.1** Each component of  $\left\langle \frac{1}{e^\rho}, \frac{2}{e^{2\rho}}, \frac{3}{e^{3\rho}}, \dots \right\rangle$

- has the sifting property:  $\int_{\rho=0}^{\rho=\infty} n e^{-n\rho} d\rho = n \frac{e^{-n\rho}}{-n} \Big|_{\rho=0}^{\rho=\infty} = 1.$
- is continuous Hyper-real function
- peaks at  $\rho = 0$  to  $n$

Therefore,

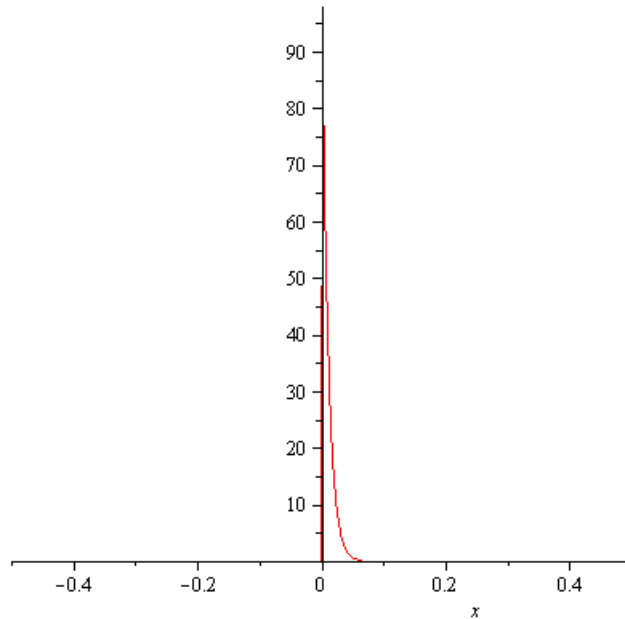
**12.2**  $\left\langle e^{-\rho}, 2e^{-2\rho}, 3e^{-3\rho}, \dots \right\rangle$  represents

*the Hyper-Real Delta Function  $\delta(\rho)$*

**12.3** For  $d\rho = \left\langle \frac{1}{n} \right\rangle$ ,  $\boxed{\delta(\rho) = \frac{1}{d\rho} e^{-\frac{\rho}{d\rho}}}$

**12.4**  $plot\left(\left\{\begin{array}{l} 0 \quad x < 0 \\ 100e^{-100x} \quad x \geq 0 \end{array}\right\}, x = -0.5 \dots 0.5\right)$  plots in Maple

the 100<sup>th</sup> component of  $\delta(r)$ ,

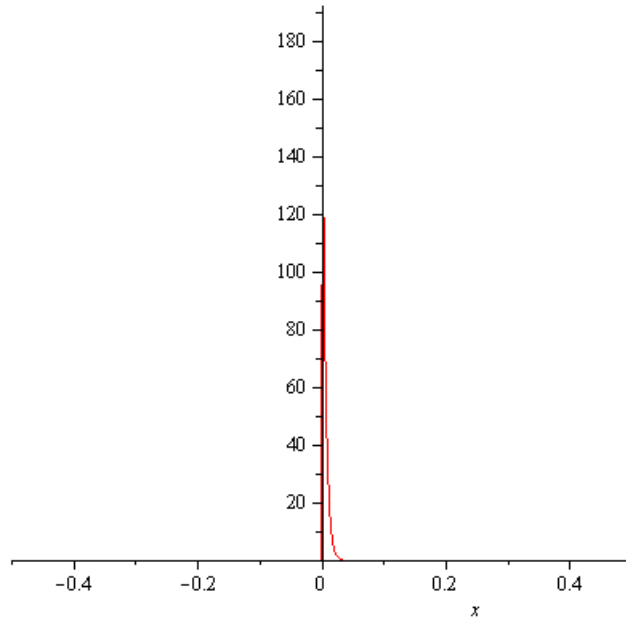


**12.5**

$plot\left(\left\{\begin{array}{l} 0 \quad x < 0 \\ 200e^{-200x} \quad x \geq 0 \end{array}\right\}, x = -0.5 \dots 0.5\right)$  plots in Maple

the 200<sup>th</sup> component of  $\delta(r)$





# 13.

## $\delta(\rho)$ and $\delta(x)\delta(y)$

$$\mathbf{13.1} \quad \delta(\rho) = \frac{1}{d\rho} \chi_{[-\frac{d\rho}{2}, \frac{d\rho}{2}]}(\rho), \quad \rho \geq 0$$

$$\mathbf{13.2} \quad \delta(\phi) = \frac{1}{d\phi} \chi_{[-\frac{d\phi}{2}, \frac{d\phi}{2}]}(\phi), \quad 0 \leq \phi \leq 2\pi$$

Transforming from Polar to Cartesian Coordinates,  $x = \rho \cos \phi$   
 $iy = i\rho \sin \phi$

$$\mathbf{13.3} \quad \delta(\rho)\delta(\phi) = \rho\delta(x)\delta(y)$$

Proof:

$$\delta(\rho) \underbrace{\delta(i\phi)}_{\frac{1}{i}\delta(\phi)} = \delta(x) \underbrace{\delta(iy)}_{\frac{1}{i}\delta(y)} \left| \frac{\partial(x, iy)}{\partial(\rho, i\phi)} \right|$$

$$\delta(\rho)\delta(\phi) = \delta(x)\delta(y) \underbrace{\begin{vmatrix} \cos \phi & i \sin \phi \\ i\rho \sin \phi & \rho \cos \phi \end{vmatrix}}_{\rho}. \quad \square$$

Integrating over  $\phi$ ,

$$\mathbf{13.4} \quad \boxed{\delta(\rho) = 2\pi\rho\delta(x)\delta(y)}$$

**Proof:** 
$$\delta(\rho) \underbrace{\int_{\phi=0}^{\phi=2\pi} \delta(\phi) d\phi}_1 = \rho \delta(x) \delta(y) \underbrace{\int_{\phi=0}^{\phi=2\pi} d\phi}_{2\pi} . \square$$

# 14.

## Cartesian Representation of $\delta(z)$

$$14.1 \quad \delta(z) = \frac{1}{i} e^{-i \operatorname{Arg}(z)} \rho \left( \frac{1}{2\pi} \int_{\omega_x=-\infty}^{\omega_x=\infty} e^{i\omega_x x} d\omega_x \right) \left( \frac{1}{2\pi} \int_{\omega_y=-\infty}^{\omega_y=\infty} e^{i\omega_y y} d\omega_y \right)$$

*Proof:* By 12.1,

$$\delta(z) = \frac{1}{2\pi i} e^{-i \operatorname{Arg}(z)} \delta(\rho)$$

By 13.4,

$$\begin{aligned} &= \frac{1}{2\pi i} e^{-i \operatorname{Arg}(z)} 2\pi \rho \delta(x) \delta(y) \\ &= \frac{1}{i} e^{-i \operatorname{Arg}(z)} \rho \delta(x) \delta(y) \end{aligned}$$

By 4.10,

$$= \frac{1}{i} e^{-i \operatorname{Arg}(z)} \rho \left( \int_{\nu_x=-\infty}^{\nu_x=\infty} e^{2\pi i \nu_x x} d\nu_x \right) \left( \int_{\nu_y=-\infty}^{\nu_y=\infty} e^{2\pi i \nu_y y} d\nu_y \right),$$

Denoting  $\omega_x = 2\pi \nu_x$ ,  
 $\omega_y = 2\pi \nu_y$ ,

$$= \frac{1}{i} e^{-i \operatorname{Arg}(z)} \rho \left( \frac{1}{2\pi} \int_{\omega_x=-\infty}^{\omega_x=\infty} e^{i\omega_x x} d\omega_x \right) \left( \frac{1}{2\pi} \int_{\omega_y=-\infty}^{\omega_y=\infty} e^{i\omega_y y} d\omega_y \right). \square$$

# 15.

## Bessel Integral Representation of

$\delta(z)$

$$\mathbf{15.1} \quad \delta(x)\delta(y) = 2\pi \int_{v=0}^{v=\infty} J_0(2\pi v\rho) v dv, \quad v = |\nu_x + i\nu_y|$$

$$= \frac{1}{2\pi} \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho) \Omega d\Omega, \quad \Omega = |\omega_x + i\omega_y|$$

= Inverse 2D-Bessel-Fourier Transform of 1.

$$\omega = 2\pi\nu, \quad \omega = \omega_x + i\omega_y, \quad \nu = \nu_x + i\nu_y$$

*Proof:* By 4.2,

$$\delta(x)\delta(y) = \left( \int_{\nu_x=-\infty}^{\nu_x=\infty} e^{2\pi i\nu_x x} d\nu_x \right) \left( \int_{\nu_y=-\infty}^{\nu_y=\infty} e^{2\pi i\nu_y y} d\nu_y \right).$$

Substitute

$$\begin{aligned} x &= \rho \cos \phi & \nu_x &= v \cos \beta \\ y &= \rho \sin \phi & \nu_y &= v \sin \beta \\ \rho &= |x + iy| & v &= |\nu_x + i\nu_y| \end{aligned}$$

Then,

$$\nu_x x + \nu_y y = v\rho(\cos \beta \cos \theta + \sin \beta \sin \theta) = v\rho \cos(\beta - \theta).$$

Integrating with respect to  $v$ , and  $\beta$ ,

$$= \int_{v=0}^{v=\infty} \left( \int_{\beta=0}^{\beta=2\pi} e^{2\pi i \nu \rho \cos(\beta-\theta)} d\beta \right) v dv$$

Denoting  $\alpha = \beta - \phi$ ,

$$\int_{\beta=0}^{\beta=2\pi} e^{2\pi i \nu \rho \cos(\beta-\phi)} d\beta = \int_{\alpha=-\phi}^{\alpha=2\pi-\phi} e^{2\pi i \nu \rho \cos \alpha} d\alpha.$$

Since  $e^{2\pi i \nu \rho \cos \alpha}$  is periodic with period  $2\pi$ ,

$$= \int_{\alpha=0}^{\alpha=2\pi} e^{2\pi i \nu \rho \cos \alpha} d\alpha,$$

$$= \int_{\alpha=0}^{\alpha=2\pi} \left( 1 + \frac{2\pi i \nu \rho \cos \alpha}{1} + \frac{(2\pi i \nu \rho \cos \alpha)^2}{2!} + \frac{(2\pi i \nu \rho \cos \alpha)^3}{3!} + \dots \right) d\alpha.$$

The integrals of the odd powers vanish, and we have

$$= 2\pi - \frac{(2\pi\nu\rho)^2}{2^2} 2\pi + \frac{(2\pi\nu\rho)^4}{2^2 \cdot 4^2} 2\pi - \frac{(2\pi\nu\rho)^6}{2^2 \cdot 4^2 \cdot 6^2} 2\pi + \dots$$

$$= 2\pi J_0(2\pi\nu\rho).$$

Therefore,

$$\delta(x)\delta(y) = 2\pi \int_{v=0}^{v=\infty} J_0(2\pi\nu\rho) v dv, \quad v = |\nu_x + i\nu_y|$$

Put  $\omega = 2\pi\nu$ ,  $\omega = \omega_x + i\omega_y$ ,  $\nu = \nu_x + i\nu_y$

$$= \frac{1}{2\pi} \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho) \Omega d\Omega, \quad \Omega = |\omega_x + i\omega_y|. \square$$

Thus,

$$\mathbf{15.2} \quad \boxed{\delta(\rho) = \rho \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho)\Omega d\Omega}$$

*Proof:* By 13.4,

$$\delta(\rho) = 2\pi\rho\delta(x)\delta(y)$$

By 15.1,

$$= 2\pi\rho \frac{1}{2\pi} \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho)\Omega d\Omega . \square$$

Hence,

$$\mathbf{15.3} \quad \boxed{\delta(z) = \frac{1}{2\pi i} e^{-i \text{Arg}(z)} \rho \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho)\Omega d\Omega}$$

*Proof:* By 11.1,

$$\delta(z) = \frac{1}{2\pi i} e^{-i \text{Arg}(z)} \delta(\rho)$$

By 15.2,

$$= \frac{1}{2\pi i} e^{-i \text{Arg}(z)} \rho \int_{\Omega=0}^{\Omega=\infty} J_0(\Omega\rho)\Omega d\Omega . \square$$

## 16.

### Primitive and Derivatives of $\delta(z)$

The Hyper-Complex Log Function is the primitive of  $\delta(z)$ , on the Hyper-Complex Plane

$$16.1 \quad \delta(\zeta - z) = \frac{d}{dz} \frac{1}{2\pi i} (\text{Log}(\zeta - z)) \mathcal{X}_{\{|\zeta - z| \leq dr\}}(\zeta)$$

$$16.2 \quad \frac{d}{dz} \delta(\zeta - z) = \frac{1}{2\pi i} \frac{1}{(\zeta - z)^2} \mathcal{X}_{\{|\zeta - z| \leq dr\}}(z)$$

That is,

- on the disk  $|\zeta - z| \leq dr$ ,  $\frac{d}{dz} \delta(\zeta - z) = \frac{1}{2\pi i} \frac{1}{(dr)^2} e^{-2i\theta}$ .
- off the disk, in  $|\zeta - z| > dr$ ,  $\frac{d}{dz} \delta(\zeta - z) = 0$ .

$$16.3 \quad \frac{d^k}{dz^k} \delta(\zeta - z) = \frac{1}{2\pi i} \frac{k!}{(\zeta - z)^{k+1}} \mathcal{X}_{\{|z| \leq dr\}}(z)$$

That is,



- on the disk,  $|\zeta - z| \leq dr$ ,  $\frac{d^k}{dz^k} \delta(\zeta - z) = \frac{k!}{2\pi i} \frac{1}{(dr)^{k+1}} e^{-i(k+1)\theta}$ ,
- off the disk, on  $|\zeta - z| > dr$ ,  $\frac{d^k}{dz^k} \delta(\zeta - z) = 0$ .

**16.4**  $\delta(z)$  is differentiable, and integrable to any order,

*But it has no Taylor Series, it is Not analytic,*

*Hence, its integral along a closed path does not vanish:*

$$\oint_{\gamma} \delta(z) dz \neq 0$$

# 17.

## Circulation of $\delta(z)$

The Hyper-complex Delta Function is Not Analytic.

It is singular on the infinitesimal disk  $|\zeta - z| \leq dr$ .

Integrating along a path that encircles its singularity at  $\zeta = z$ , the Circulation of Delta along the infinitesimal circle is 1.

### 17.1

$$\oint_{|\zeta-z|=dr} \delta(\zeta - z) d\zeta = 1.$$

*Proof:* Put

$$\begin{aligned}\zeta - z &= (dr)e^{i\alpha} \\ d\zeta &= i(dr)e^{i\alpha} d\alpha.\end{aligned}$$

Then,

$$\frac{1}{2\pi i} \oint_{|\zeta-z|=dr} \frac{1}{\zeta - z} d\zeta = \frac{1}{2\pi i} \int_{\alpha=0}^{\alpha=2\pi} \frac{1}{(dr)e^{i\alpha}} i(dr)e^{i\alpha} d\alpha,$$

Since  $(dr)e^{i\alpha} \neq 0$ , for any infinitesimal  $dr$ , and any  $\alpha$ ,

$$= \frac{1}{2\pi i} i \int_{\alpha=0}^{\alpha=2\pi} d\alpha = 1. \square$$

# 18.

## Sifting by $\delta(\zeta - z)$ and $\frac{d}{dz}\delta(\zeta - z)$

### 18.1 Sifting by $\delta(\zeta - z)$

If  $f(z)$  is Hyper-Complex Differentiable function at  $z$

Then,

$$\oint_{|\zeta-z|=dr} f(\zeta)\delta(\zeta - z)d\zeta = f(z)$$

Proof:

Since  $f(\zeta)$  is differentiable at  $z$ , on the circle  $\zeta - z = (dr)e^{i\alpha}$ ,

$$\zeta = z + (dr)e^{i\alpha},$$

$$d\zeta = i(dr)e^{i\alpha}d\alpha$$

$$f(z + (dr)e^{i\alpha}) = f(z) + f'(z)(dr)e^{i\alpha},$$

$$\oint_{|\zeta-z|=dr} f(\zeta)\delta(\zeta - z)d\zeta =$$

$$= f(z) \underbrace{\oint_{|\zeta-z|=dr} \delta(\zeta - z)d\zeta}_1 + f'(z)dr \oint_{|\zeta-z|=dr} e^{i\alpha} \frac{1}{2\pi i} \underbrace{\frac{1}{\zeta - z}}_{\frac{1}{dr}e^{-i\alpha}} \underbrace{d\zeta}_{i(dr)e^{i\alpha}d\alpha}$$

$$= f(z) + \frac{1}{2\pi i} f'(z)(dr) \int_{\alpha=0}^{\alpha=2\pi} e^{i\alpha} \frac{1}{(dr)e^{i\alpha}} i(dr)e^{i\alpha} d\alpha$$

$$\begin{aligned}
 &= f(z) + \frac{1}{2\pi i} f'(z)(dr) \underbrace{\int_{\alpha=0}^{\alpha=2\pi} e^{i\alpha} d\alpha}_0 \\
 &= f(z). \square
 \end{aligned}$$

**18.2**

$$\boxed{\frac{d}{dz} f(z) = \oint_{|\zeta-z|=dr} f(\zeta) \frac{d}{dz} \delta(\zeta - z) d\zeta}$$

Proof:

$$\begin{aligned}
 \frac{d}{dz} f(z) &= \frac{1}{2\pi i} \oint_{|\zeta-z|=dr} \frac{f(\zeta)}{(\zeta - z)^2} d\zeta \\
 &= \oint_{|\zeta-z|=dr} f(\zeta) \frac{d}{dz} \delta(\zeta - z) d\zeta . \square
 \end{aligned}$$

**18.3**

$$\boxed{\frac{d^k}{dz^k} f(z) = \oint_{|\zeta-z|=dr} f(\zeta) \frac{d^k}{dz^k} \delta(\zeta - z) d\zeta}$$

Proof:

$$\begin{aligned}
 \frac{d}{dz} f(z) &= \frac{k!}{2\pi i} \oint_{|\zeta-z|=dr} \frac{f(\zeta)}{(\zeta - z)^{k+1}} d\zeta \\
 &= \oint_{|\zeta-z|=dr} f(\zeta) \frac{d^k}{dz^k} \delta(\zeta - z) d\zeta . \square
 \end{aligned}$$

# 19.

## $\delta(\zeta - z)$ and the Cauchy Integral Formula

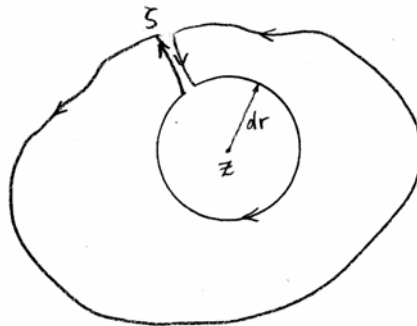
### 19.1 The Cauchy Integral Formula is Sifting by Delta

If  $f(z)$  is Hyper-Complex Differentiable function on a Hyper-Complex Simply-Connected Domain  $D$ .

Then, for any loop  $\gamma$ , and any point  $z$  in its interior

$$f(z) = \oint_{|z-\zeta|=dr} f(\zeta)\delta(\zeta - z)d\zeta = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(\zeta)}{\zeta - z} d\zeta$$

Proof: The Hyper-Complex function  $\frac{f(\zeta)}{\zeta - z}$  is Differentiable on the Hyper-Complex Simply-Connected domain  $D$ , and on a path that includes  $\gamma$  and an infinitesimal circle about  $z$ .



Then, the integrals on the lines between  $\gamma$  and the circle have opposite signs and cancel each other.

The integral over the circle has a negative sign because its direction is clockwise, and by Cauchy Integral Theorem,

$$\oint_{\gamma} \frac{f(\zeta)}{\zeta - z} d\zeta - \oint_{|\zeta-z|=dr} \frac{f(\zeta)}{\zeta - z} d\zeta = 0.$$

Therefore,

$$\begin{aligned} \oint_{\gamma} \frac{f(\zeta)}{\zeta - z} d\zeta &= \oint_{|\zeta-z|=dr} \frac{f(\zeta)}{\zeta - z} d\zeta \\ &= 2\pi i \underbrace{\oint_{|\zeta-z|=dr} f(\zeta) \frac{1}{2\pi i} \frac{1}{\zeta - z} d\zeta}_{f(z)}. \square \end{aligned}$$

## 20.

# $\delta(\zeta - z)$ and the Residue of a Laurent Expansion

### 20.1 Laurent Expansion of a Singular $f(z)$

If  $f(z)$  is Hyper-Complex Differentiable function on a Hyper-Complex disk  $0 < |z - z_0| < r$

Then,  $f(z) = \dots + a_{-3}(z - z_0)^{-3} + a_{-2}(z - z_0)^{-2} + a_{-1}(z - z_0)^{-1} +$   
 $+ a_0 + a_1(z - z_0) + a_2(z - z_0)^2 + \dots$

where for any loop  $\gamma$ , and for any point  $z \neq z_0$  in its interior

$$a_{-k} = \frac{1}{2\pi i} \oint_{\gamma} f(\zeta)(\zeta - z_0)^{k-1} d\zeta, \quad k = 1, 2, \dots,$$

$$a_{-3} = \frac{1}{2\pi i} \oint_{\gamma} f(\zeta)(\zeta - z_0)^2 d\zeta$$

$$a_{-2} = \frac{1}{2\pi i} \oint_{\gamma} f(\zeta)(\zeta - z_0) d\zeta,$$

$$a_{-1} = \frac{1}{2\pi i} \oint_{\gamma} f(\zeta) d\zeta,$$

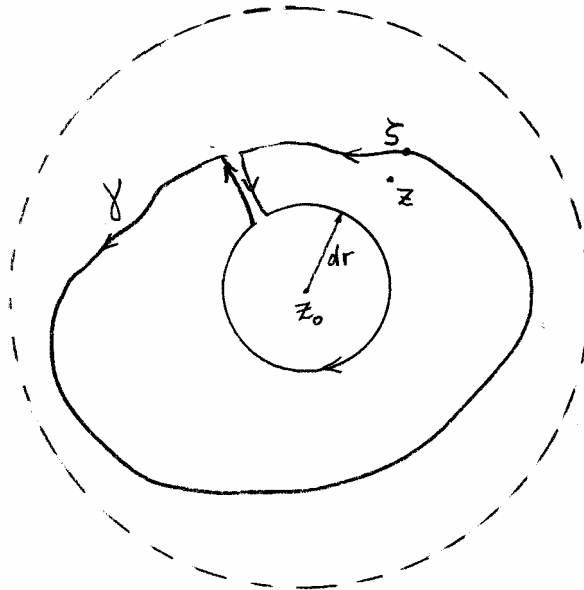
$$a_0 = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(\zeta)}{\zeta - z_0} d\zeta,$$

$$a_1 = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(\zeta)}{(\zeta - z_0)^2} d\zeta,$$

$$a_2 = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(\zeta)}{(\zeta - z_0)^3} d\zeta$$

$$a_k = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(\zeta)}{(\zeta - z_0)^{k+1}} d\zeta, \quad k = 0, 1, 2, \dots$$

**Proof:** The Hyper-Complex Differentiable  $f(z)$  satisfies Cauchy Integral Formula in the Hyper-Complex domain  $D$ , bounded by a path that includes  $\gamma$  and an infinitesimal circle about  $z_0$





Then, the integrals on the lines between  $\gamma$  and the circle have opposite signs and cancel each other.

The integral over the circle has a negative sign because its direction is clockwise, and by Cauchy Integral Formula,

$$f(z) = \frac{1}{2\pi i} \left( \oint_{\gamma} \frac{f(\zeta)}{\zeta - z} d\zeta - \oint_{|\zeta - z| = dr} \frac{f(\zeta)}{\zeta - z} d\zeta \right).$$

For  $\zeta$  along  $\gamma$ ,

$$\frac{1}{\zeta - z} = \frac{1}{\zeta - z_0} \frac{1}{1 - \frac{z - z_0}{\zeta - z_0}} = \frac{1}{\zeta - z_0} \left( 1 + \frac{z - z_0}{\zeta - z_0} + \left( \frac{z - z_0}{\zeta - z_0} \right)^2 + \dots \right).$$

Then,

$$\oint_{\gamma} \frac{f(\zeta)}{\zeta - z} d\zeta = \underbrace{\oint_{\gamma} \frac{f(\zeta)}{\zeta - z_0} d\zeta}_{2\pi i a_0(z_0)} + \underbrace{\oint_{\gamma} \frac{f(\zeta)}{(\zeta - z_0)^2} d\zeta}_{2\pi i a_1(z_0)} (z - z_0) + \underbrace{\oint_{\gamma} \frac{f(\zeta)}{(\zeta - z_0)^3} d\zeta}_{2\pi i a_2(z_0)} (z - z_0)^2 + \dots$$

For  $\zeta$  along the circle  $|\zeta - z| = dr$ ,

$$-\frac{1}{\zeta - z} = \frac{1}{z - z_0} \frac{1}{1 - \frac{\zeta - z_0}{z - z_0}} = \frac{1}{z - z_0} \left( 1 + \frac{\zeta - z_0}{z - z_0} + \left( \frac{\zeta - z_0}{z - z_0} \right)^2 + \dots \right).$$

Then,

$$\oint_{|\zeta - z| = dr} \frac{-f(\zeta)}{\zeta - z} d\zeta = \underbrace{\oint_{|\zeta - z| = dr} f(\zeta) d\zeta}_{2\pi i a_{-1}(z_0)} \frac{1}{z - z_0} + \underbrace{\oint_{|\zeta - z| = dr} \frac{f(\zeta)}{\zeta - z_0} d\zeta}_{2\pi i a_{-2}(z_0)} \frac{1}{(z - z_0)^2} + \underbrace{\oint_{|\zeta - z| = dr} \frac{f(\zeta)}{(\zeta - z_0)^2} d\zeta}_{2\pi i a_{-3}(z_0)} \frac{1}{(z - z_0)^3} + \dots$$

By the Cauchy Integral Theorem the integrals of  $a_{-1}, a_{-2}, a_{-3}, \dots$  can be taken along  $\gamma$ .  $\square$

**20.2  $\delta(\zeta - z)$  and the Residue of a Laurent Expansion**

If  $f(z)$  is Hyper-Complex function on a Hyper-Complex disk

$0 < |z - z_0| < r$  so that

$$f(z) = \dots + a_{-3}(z - z_0)^{-3} + a_{-2}(z - z_0)^{-2} + a_{-1}(z - z_0)^{-1} + \\ + a_0 + a_1(z - z_0) + a_2(z - z_0)^2 + \dots$$

Then, for any loop  $\gamma$ , around  $z_0$

$$\oint_{\gamma} (z - z_0)^k dz = \begin{cases} 0, & k \neq -1 \\ 2\pi i \oint_{|z-z_0|=dr} \delta(z - z_0) dz = 2\pi i, & k = -1 \end{cases}$$

Hence,

$$\frac{1}{2\pi i} \oint_{\gamma} f(\zeta) d\zeta = a_{-1},$$

Proof:

For any integer  $k \neq -1$ ,

$$\oint_{\gamma} (z - z_0)^k dz = \frac{1}{k + 1} (z - z_0)^{k+1} \Big|_{z=\alpha}^{z=\alpha} = 0$$

For  $k = -1$ ,

$$\oint_{\gamma} (z - z_0)^{-1} dz = 2\pi i \oint_{|z-z_0|=dr} \frac{1}{2\pi i} \frac{1}{z - z_0} dz = 2\pi i . \square$$

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